

# On Matérn-type Gaussian processes

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## Matérn correlations, context

Matérn correlations are very popular, their uses date at least to Matheron 1960, maybe Krige ... A Matérn isotropic zero-mean stationary Gaussian field model for (possibly noisy) observations at sites  $\vec{s}_i \in \mathbb{R}^d, i = 1, \dots, n$ , with (known) regularity (an index  $\nu > 0$ ), is determined by,  $\|\cdot\|$  being the Euclidean norm in  $\mathbb{R}^d$ :

$$\mathbb{E}(Z(\vec{s})Z(\vec{t})) = b \rho_{\nu,\theta}(\|\vec{s} - \vec{t}\|), \quad \mathbb{E}(Z(\vec{s})^2) \equiv b > 0$$

where

$$\rho_{\nu,\theta}(x) := \frac{(\theta x)^\nu}{\Gamma(\nu)2^{\nu-1}} K_\nu(\theta x), \quad x > 0, \quad \theta > 0,$$

where  $K_\nu$  is the Bessel function of the 2nd kind of order  $\nu > 0$ .

- $\nu = 1/2$ :  $\rho_{\nu,\theta}(x) = e^{-\theta x}$  (“Exponential” case. In 1D: O.U.)
- $\nu = 3/2$ :  $\rho_{\nu,\theta}(x) = (1 + \theta x)e^{-\theta x}$
- $\nu = +\infty$ :  $\lim_{\nu \rightarrow \infty} \rho_{\nu, \frac{2\nu^2}{\ell}}(x) = e^{-\frac{x^2}{\ell^2}}$  (“squared exponential”)

## Spectral densities

Models more succinctly defined by their family of spectral densities over  $(-\infty, +\infty)^d$ :

$$f_{\nu, b, \theta}^*(\vec{\omega}) = b g_{\nu, \theta}^*(\vec{\omega}), \quad \text{with } g_{\nu, \theta}^*(\vec{\omega}) := \frac{C(d, \nu) \theta^{2\nu}}{(\theta^2 + \|\vec{\omega}\|^2)^{\nu + \frac{d}{2}}}$$

where the constant  $C(d, \nu)$  is chosen s.t.  $\int_{\mathbf{R}^d} g_{\nu, \theta}^*(\vec{\omega}) d\vec{\omega} = 1$

Here **one mainly considers gridded data sites, observed on the hypercube (i.e. data sites are  $\delta\{1, \dots, n\}^{\otimes d}$ )**

From the well known aliasing formula, the spectral density on  $(-\pi, \pi]^d$  of  $Z_\delta$  is

$$f_{\nu, \theta}^\delta(\vec{\lambda}) = b g_{\nu, \theta}^\delta(\vec{\lambda}) \quad \text{with } g_{\nu, \theta}^\delta(\vec{\lambda}) := \frac{1}{\delta^d} \sum_{\vec{k} \in \mathbf{Z}^d} g_{\nu, \theta}^* \left( \frac{\vec{\lambda} + 2\vec{k}\pi}{\delta} \right).$$

## Spectral densities, and quasi-Matérn variant

Simple expressions for  $g_{\nu,\theta}^{\delta}(\vec{\lambda})$  are available only for  $d = 1$ , when  $\nu - 1/2$  is a small integer, says 0, 1 or 2 : they then coincide with particular ARMA spectral densities with a *constrained* vector of parameters.

The quasi-Matérn spectral density model one considers here, for a process defined on  $\mathbb{Z}^d$ , is defined by  $f_{\mathcal{Q},\nu,\beta}(\vec{\lambda}) = b g_{\mathcal{Q},\nu,\beta}(\vec{\lambda})$  with

$$g_{\mathcal{Q},\nu,\beta}(\vec{\lambda}) := C_{\mathcal{Q}}(d, \nu) \left( d + \beta - \sum_{k=1}^d \cos(\lambda_k) \right)^{-\nu-d/2}, \quad \beta > 0,$$

$C_{\mathcal{Q}}(d, \nu)$  being again chosen s.t.  $\int_{[\pi, \pi]^d} g_{\mathcal{Q},\nu,\beta}(\vec{\lambda}) d\vec{\lambda} = 1$

- the case  $d = 2, \nu = 1$  gives the SAR model of order 1

Motivation : As  $\delta \downarrow 0$ ,  $g_{\nu,\theta}^{\delta}(\vec{\lambda}) \simeq g_{\mathcal{Q},\nu,\beta}(\delta\vec{\lambda})$  (with  $\beta = (\delta\theta)^2/2$ ).

## A not so “well-known” expression

Let  $\vartheta_3$  denote the Jacobi 3rd theta function. Recall that

$$\sqrt{\frac{\pi}{\alpha^2}} \sum_{j=-\infty}^{\infty} e^{-\frac{(2\pi j + \lambda)^2}{4\alpha^2}} = \vartheta_3\left(\frac{\lambda}{2}, e^{-\alpha^2}\right)$$

### Proposition 1

With  $\alpha = \delta\theta$  and  $w_\nu(\kappa) := \frac{e^{-\frac{1}{4\kappa}} \kappa^{-\nu-1}}{2^{2\nu} \Gamma(\nu)}$  (it is a pdf on  $[0, \infty[$ ), it holds

$$g_{\nu, \theta}^\delta(\lambda_1, \dots, \lambda_d) = \int_0^\infty g_{\infty, \alpha\phi}(\lambda_1, \dots, \lambda_d) \times 2\phi w_\nu(\phi^2) d\phi$$

where

$$g_{\infty, \alpha}(\vec{\lambda}) = \frac{1}{(2\pi)^d} \prod_{k=1}^d \vartheta_3\left(\frac{\lambda_k}{2}, e^{-\alpha^2}\right).$$

## Two “monotonicity” results (or conjectures)

Let  $d = 1, 2, \dots$  and  $\nu > 0$  fixed. In the following,  $g_\alpha(\cdot)$  stands either for  $g_{\nu, \alpha}^{\delta=1}(\cdot)$  (which coincides with  $g_{\nu, \frac{\alpha}{\delta}}^\delta(\cdot)$ ) or for  $g_{\mathcal{Q}, \nu, \alpha}(\cdot)$ .

### Monotonicity of the “entropy-rate”

$$\alpha \in \mathbb{R}_+ \longrightarrow \int_{[-\pi, \pi]^d} \log \left( g_\alpha(\vec{\lambda}) \right) d\vec{\lambda}$$

is strictly increasing.

### ‘Unimodality’ property. Let $\alpha_0$ fixed ( $> 0$ )

$$\alpha \in \mathbb{R}_+ \longrightarrow \frac{1}{(2\pi)^d} \int_{[-\pi, \pi]^d} \frac{g_{\alpha_0}(\vec{\lambda})}{g_\alpha(\vec{\lambda})} d\vec{\lambda} - 1 =: F(\alpha_0, \alpha)$$

has a unique finite root  $\alpha = \alpha_0$ .

## Two “monotonicity” results (or conjectures)

Their proofs will be (hopefully) published soon.

- They are easy for the quasi-Matérn spectral densities.
- For pure Matérn spectral densities, the proofs are relatively easy for  $d = 1$ .

## Consequence of this “entropy monotonicity” property

As is well known, in 1D, this entropy rate is proportional to the prediction error variance of  $Z(t = 0)$  based on the past (Szegő's formula).

Generalization of the Szegő's formula is also known in 2D when the “past” is a half-plane.

This “entropy” monotonicity thus indicates that the ability to extrapolate at unobserved sites (in geostatistic it's called the **ability to interpolate**), of such processes, whose the underlying true range parameter is  $\theta$ , is strengthened when  $\theta$  increases.



## Consequences of this “Unimodality” property

The so-called “(Gaussian) energy-variance matching” (GE-EV) method for estimating  $\alpha_0$  is an empirical version of the above equation  $F(\alpha_0, \alpha) = 0$  in  $\alpha$ .

Recall that the (maybe simplest version of) **GE–EV estimating equation is defined by replacing, in this integrated ratio,  $g_{\alpha_0}(\cdot)$  by a tapered periodogram normalized by  $\frac{1}{n^d} \sum_{s \in \delta\{1, \dots, n\}^{\otimes d}} Z(s)^2$ .**

The previously published consistency for GE–EV was only a weak result (“there exists a sequence of roots ...”) and required the assumption :  $\alpha_0$  (or  $\delta$ ) must be small..

**This “unimodality” property implies a global consistency result, even without this assumption.**

It also entails some “robustness” properties for simple algorithms for computing the GE–EV estimate of  $\alpha_0$

## A third “monotonicity” result (or conjecture...)

Let  $h_\alpha(\vec{\lambda}) = \frac{\partial}{\partial \alpha} \log(g_\alpha(\vec{\lambda}))$ . Define the “coefficient of variation” of  $f$  (where “ $\int$ ” denotes integrals over  $[-\pi, \pi]^d$ ) by:

$$\text{CV}(f) := \left\{ \frac{1}{(2\pi)^d} \int \left( \left| f - \left( \frac{1}{(2\pi)^d} \int f \right) \right|^2 \right) \right\} / \left( \frac{1}{(2\pi)^d} \int f \right)^2$$

Monotonicity property of the “inefficiency”

$$\alpha \in \mathbb{R}_+ \longrightarrow \text{CV}(h_\alpha) \text{CV}(g_\alpha)$$

is a decreasing function.

## Consequences of this “CVs product monotonicity”

Assume the process is Gaussian Matérn (or Gaussian quasi-Matérn) with sdf  $b_0 g_{\alpha_0}$  (recall that  $\alpha_0 = \delta\theta_0$  for pure Matérn).

inefficiency $_{\nu, \alpha_0} :=$  “asymptotic variance” (a.v.) of the GE–EV estimator divided by the a.v. of the maximum likelihood estimator .

Recall

Theorem (slight extension of Girard 2016).

For  $d = 1$

$$\text{inefficiency}_{\nu, \alpha_0} = \text{CV}(h_{\alpha_0}) \text{CV}(g_{\alpha_0}) \rightarrow \text{ineff}_1(\nu) \quad \text{as } \alpha_0 \downarrow 0,$$

$$\text{where } \text{ineff}_1(\nu) := \frac{\sqrt{\pi}}{2} \left( \frac{\Gamma(\nu+3/2)}{\Gamma(\nu+1)} \right)^2 \frac{\Gamma(2\nu+1/2)}{\Gamma(2\nu+1)}.$$

The above approximates values of inefficiency $_{\nu, \alpha_0}$  for the cases when the true  $\alpha_0$  is small are satisfactory since they are close to 1 for typical values of  $\nu$ . For instance

$$\text{ineff}_1(1/2) = 1, \quad \text{ineff}_1(3/2) = 1.054^2, \quad \text{ineff}_1(5/2) = 1.122^2$$

Conjecture (extension of the near-efficiency result of Girard 2016 to any  $\alpha_0$  (or to any  $\delta$  in the case of pure Matérn)).

$$\forall \alpha_0 > 0, \quad \text{inefficiency}_{\nu, \alpha_0} \leq \text{ineff}_d(\nu)$$

where  $\text{ineff}_d(\nu) := \lim_{\alpha \downarrow 0} \text{CV}(h_\alpha) \text{CV}(g_\alpha)$ .

## Theorem.

For  $d = 1$ , the above conjecture holds true for  $\nu = 1/2$  a “small” integers; and thus, in particular, for any true  $\alpha_0$

- inefficiency $_{\nu, \alpha_0} \equiv 1$  in the case  $\nu = 1/2$ ,
- inefficiency $_{\nu, \alpha_0} \leq 10/9$  in the case  $\nu = 3/2$ ,
- inefficiency $_{\nu, \alpha_0} \leq 63/50$  in the case  $\nu = 5/2$ ,
- inefficiency $_{\nu, \alpha_0} \leq \frac{1716}{1225} \simeq 1.18356^2$  in the case  $\nu = 7/2$ .

In 2D, a numerical verification for quasi-Matérn <sub>$\frac{1}{2}$</sub>

$$\nu = \frac{1}{2}$$

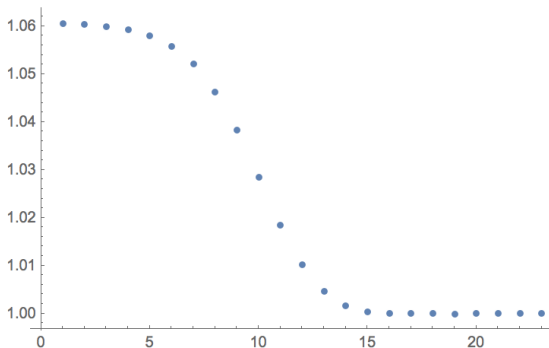


Figure 1 :  $\alpha_0 \rightarrow \sqrt{(\text{inefficiency}_{\nu, \alpha_0})}$

In 2D, for quasi-Matérn $_{\frac{3}{2}}$

$$\nu = \frac{3}{2}$$

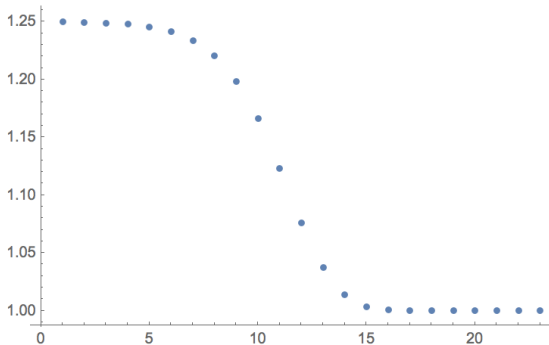


Figure 2 :  $\alpha_0 \rightarrow \sqrt{(\text{inefficiency}_{\nu, \alpha_0})}$

## A finite size analog of “monotonicity of the entropy-rate”

Let  $d = 1, 2, \dots$  and  $\nu > 0$  fixed. Let  $S = \{\vec{s}_i, i = 1, \dots, n\}$  fixed. In the following,  $R_{\nu, \theta}$  denotes the  $n \times n$  matrix with  $(i, j)$ th term  $:= \rho_{\nu, \theta}(\|\vec{s}_i - \vec{s}_j\|)$

### Monotonicity of the determinant

$$\theta \in \mathbb{R}_+ \longrightarrow \det(R_{\nu, \theta})$$

is strictly increasing.

**Theorem.** If  $\nu = \frac{1}{2}$  then

this Monotonicity of the determinant holds true.

**Proof.** It stems on  $R_{\nu, \theta_1 + \theta_2} \equiv R_{\nu, \theta_1} \circ R_{\nu, \theta_2}$  when  $\nu = 1/2$ , and a known determinant inequality for Hadamard products (Oppenheim 1930).



## Consequence of this determinant's monotonicity

Let  $F_n(\theta_0, \theta) := \frac{1}{n} \text{trace} \left( R_{\nu, \theta}^{-1} R_{\nu, \theta_0} \right)$ .

It is known that

$$F_n(\theta_0, \theta) \leq \left( \det \left( R_{\nu, \theta}^{-1} R_{\nu, \theta_0} \right) \right)^{\frac{1}{n}}$$

**Theorem.** If  $\nu = \frac{1}{2}$  then

the function

$$\theta \longrightarrow F_n(\theta_0, \theta) - 1$$

has no root in  $\theta \in ]0, \theta_0[$ .

**NB:** The GE-EV method for estimating  $\theta_0$  is an empirical version of the above equation  $F_n(\theta_0, \theta) = 1$  in  $\theta$ .

## Some References

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**THANK YOU FOR YOUR ATTENTION**